Traces of Thermalization at RHIC

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onset of thermal equilibration – common centrality dependence of $\langle p_t \rangle$, p_t and charge dynamic fluctuations

- I. Fluctuations from nonequilibrium two-body correlations
- multiplicity and p_t fluctuation observables
- approach to equilibrium
- initial and near-equilibrium fluctuations
- II. Experiments
- ▶ PHENIX and STAR mean p_t and fluctuations
- energy dependence, net charge fluctuations

nucl-th/0308067

Dynamic Fluctuations

variance minus thermal contribution

Pruneau, Voloshin & S.G.

multiplicity
$$N$$

$$R_{AA} = \frac{\left\langle N^2 \right\rangle - \left\langle N \right\rangle^2 - \left\langle N \right\rangle}{\left\langle N \right\rangle^2}$$

mean
$$p_t$$

$$\langle \delta p_{t1} \delta p_{t2} \rangle \equiv \frac{1}{\langle N(N-1) \rangle} \left\langle \sum_{i \neq j} \delta p_{ti} \delta p_{tj} \right\rangle$$

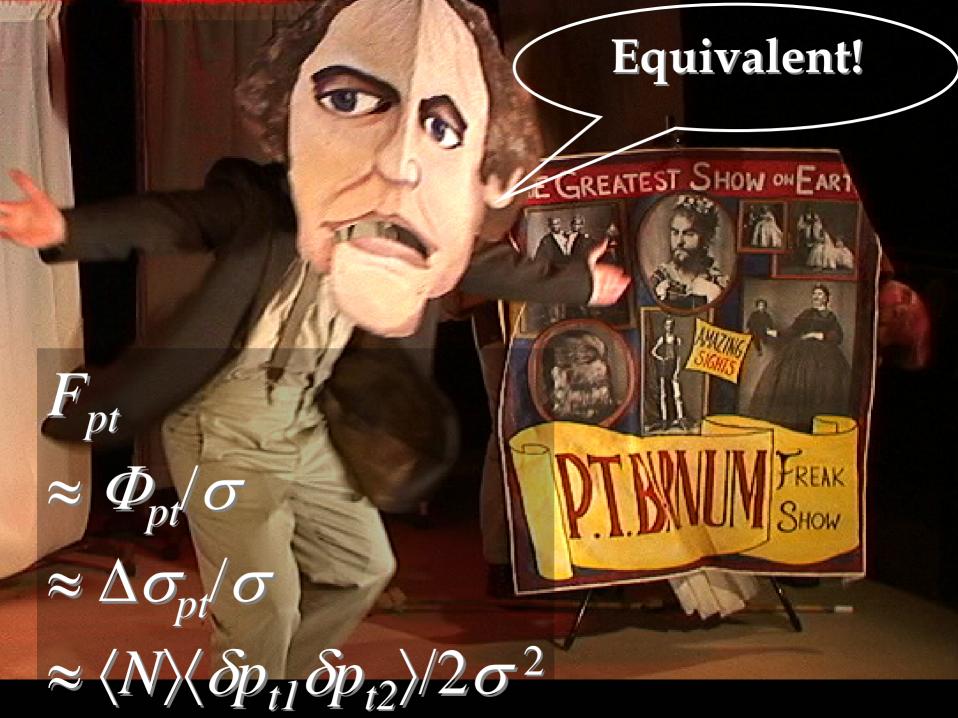
$$\delta p_{\scriptscriptstyle t} \equiv p_{\scriptscriptstyle t} - \langle p_{\scriptscriptstyle t} \rangle$$

probe two-body correlation function:

$$r(p_1, p_2) = \frac{dN}{dp_1 dp_2} - \frac{dN}{dp_1} \frac{dN}{dp_2}$$

$$R_{AA} \propto \iint r(p_1, p_2)$$

$$\langle \delta p_{t1} \delta p_{t2} \rangle \propto \iint \delta p_{t1} \delta p_{t2} \, r(p_1, p_2)$$



Time Scales

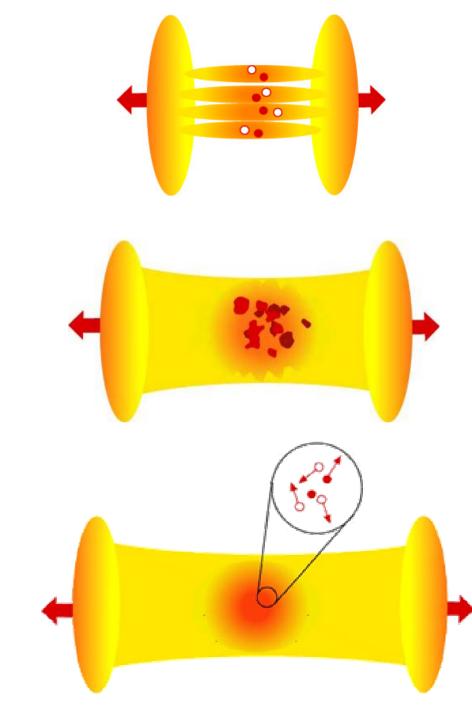
initially – string fragmentation

later – clumps, size $\sim \xi$

local thermalization; time $\sim v^{-1}$ scattering rate $v = \langle \sigma | v_{rel} \rangle n$

much later (if ever)

- flow between clumps → homogeneity
- diffusion time $t_{diff} \sim v \xi^2$



Fluctuation Sources

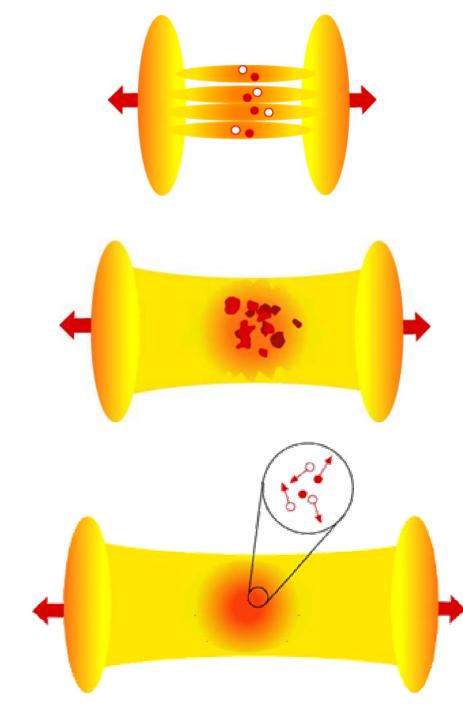
dynamic fluctutations ~ (strings)⁻¹

dynamic fluctutations ~ (clumps)⁻¹

-- larger? depends on clump size

Gazdzicki & Mrowczynski

statistical fluctuations ~ (particles)⁻¹ dynamic fluctuations **vanish**

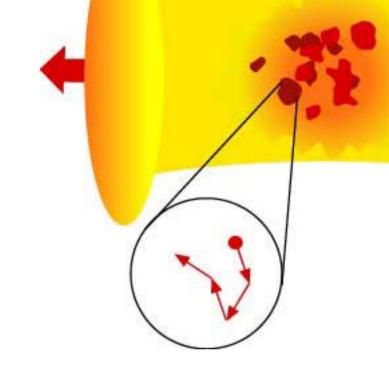


Partial Thermalization

random walk – n collisions

$$\langle p_t^2 \rangle = \langle p_t^2 \rangle_0 + \kappa (n-1)$$

collision: many walkers bouncing off each other – energy conservation limits p_t increase



Boltzmann equation
$$\Rightarrow \langle p_t^2 \rangle = \langle p_t^2 \rangle_0 S + \langle p_t^2 \rangle_e (1 - S)$$

survival probability:

$$S = e^{-\int v \, d\tau} \approx \left(\tau_0 \, / \, \tau_F\right)^{\alpha}$$

few collisions
$$\langle p_t^2 \rangle_e - \langle p_t^2 \rangle_0 \sim \kappa$$

many collisions – saturates at $\langle p_{_t}^2 \rangle$

Nonequilibrium $\langle p_t \rangle$

central collisions: more participants N

longer lifetime \Rightarrow smaller survival probability S

$$\langle p_t \rangle = \langle p_t \rangle_0 S + \langle p_t \rangle_e (1 - S)$$

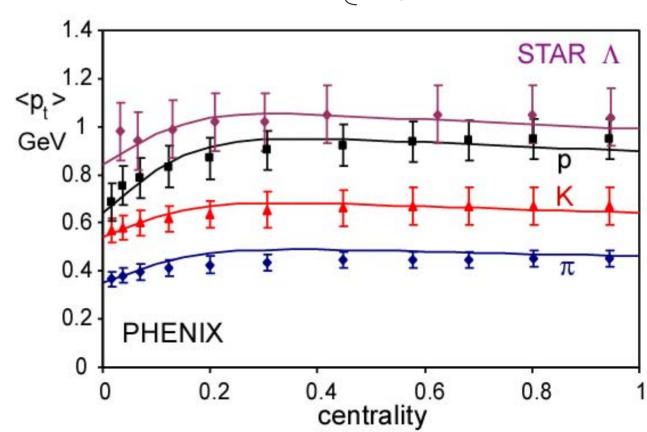
Boltzmann equation + longitudinal expansion

survival probability

$$S = \left(t_0 / t_F\right)^{\alpha}$$

equilibrium cooling $\left\langle p_{t}\right\rangle _{a}\propto\left(t_{0}\left/t_{F}\right)^{\gamma}$

centrality dependence $t_E \propto N^x$, $\alpha \propto N^y$



Nonequilibrium Dynamic Fluctuations

Boltzmann equation with Langevin noise ⇒ phase-space correlations ⇒ dynamic fluctuations

simple limits:

independent initial and near-local equilibrium correlations

$$\langle \delta p_{t1} \delta p_{t2} \rangle = \langle \delta p_{t1} \delta p_{t2} \rangle_0 S^2 + \langle \delta p_{t1} \delta p_{t2} \rangle_{le} (1 - S)^2$$

initial correlations equilibrium-like

$$\left\langle \delta p_{t1} \delta p_{t2} \right\rangle = \left\langle \delta p_{t1} \delta p_{t2} \right\rangle_{0} S^{2} + \left\langle \delta p_{t1} \delta p_{t2} \right\rangle_{le} \left(1 - S^{2} \right)$$

same centrality dependence of S

Initial and Near-Local Equilibrium Fluctuations

initial fluctuations:

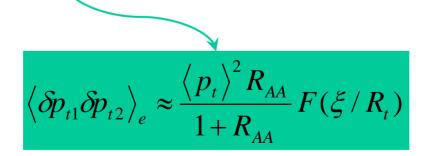
M participant nucleons

→ independent strings

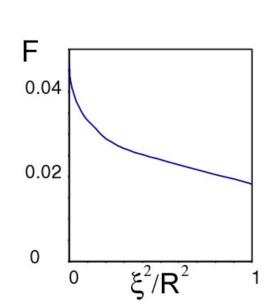
$$\langle \delta p_{t1} \delta p_{t2} \rangle \approx \frac{2 \langle \delta p_{t1} \delta p_{t2} \rangle_{pp}}{M} \left(\frac{1 + R_{pp}}{1 + R_{AA}} \right)$$

near equilibrium: correlations from clumpiness – more likely to find particles near "hot spots" \rightarrow spatial correlation function $r(x_1, x_2)$

$$\langle \delta p_{t1} \delta p_{t2} \rangle \propto \iint \delta \overline{p}_{t1}(x_1) \delta \overline{p}_{t2}(x_2) r(x_1, x_2) dx_1 dx_2$$



transverse size R_r , correlation length ξ



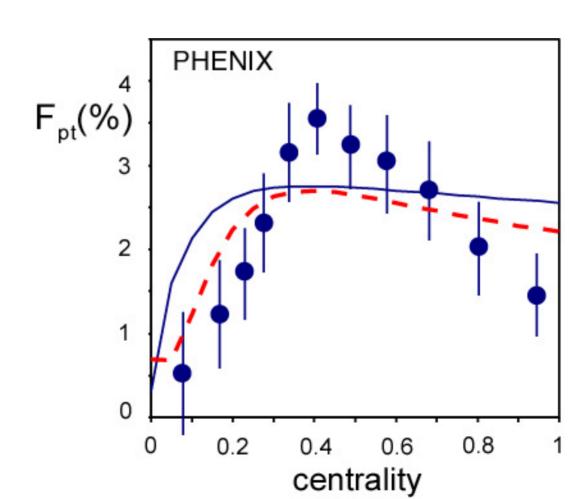
Nonequilibrium Fluctuations

$$\langle \delta p_{t1} \delta p_{t2} \rangle = \langle \delta p_{t1} \delta p_{t2} \rangle_0 S^2 + \langle \delta p_{t1} \delta p_{t2} \rangle_{le} (1 - S^2)$$

initial fluctuations – wounded nucleons

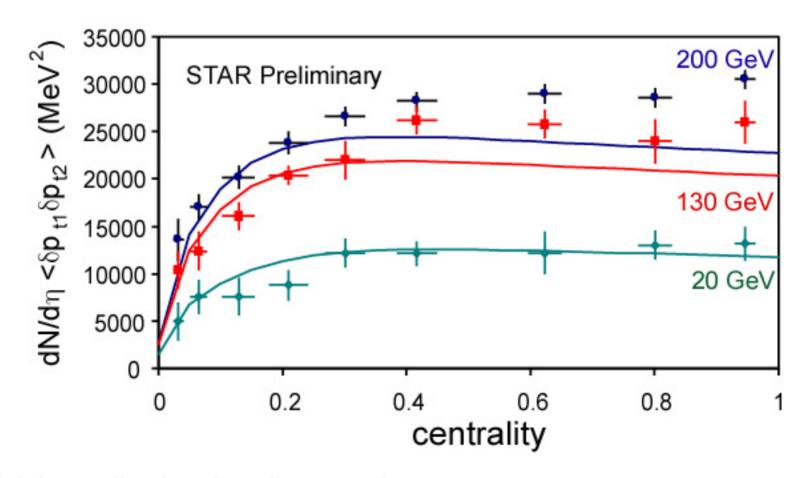
near equilibrium correlation length $\xi \sim 1$ fm; R $\propto N^{1/2}$

survival probability centrality dependence from $\langle p_t \rangle$ data



Energy Dependence

multiplicity, scattering rate $\alpha \propto dN/d\eta$, $R_{AA} \propto (dN/d\eta)^{-1}$



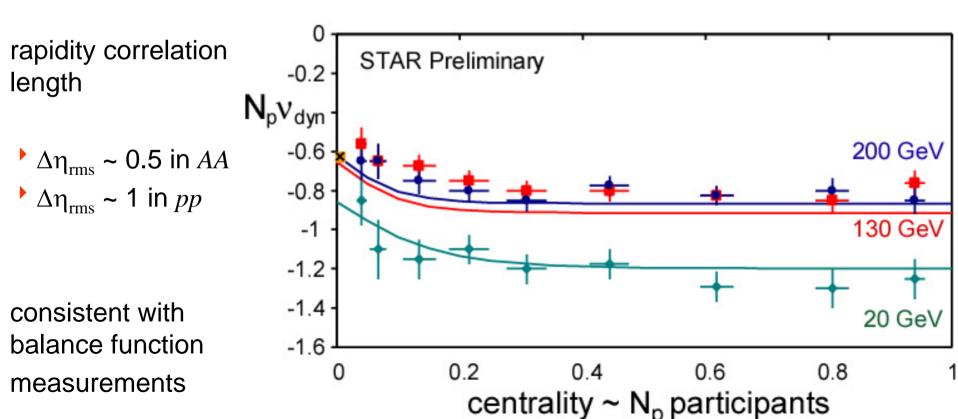
FIND: partial thermalization describes trends.

Deviation in central collisions at highest energies – jets?

Similar Charge Fluctuations

expect same charge fluctuations \rightarrow correlations $r(p_1,p_2)$

$$v_{dyn} = R_{++} + R_{--} - 2R_{+-} \longrightarrow v_0 S^2 + v_{le} (1 - S^2)$$



Fluctuations: not just for Exotic Phenomena!

$\langle p_t \rangle$, p_t and charge fluctuations

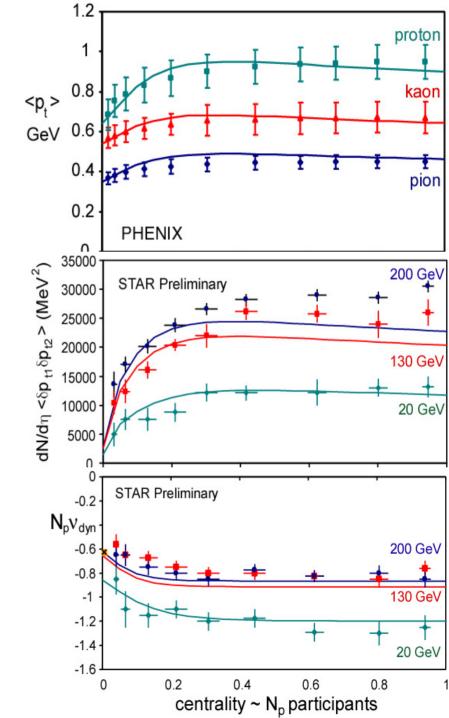
- common behavior peripheral
- but: alternative indivdual explanations

central collisions?

- cooling? PHENIX: yes, STAR: no
- jet contribution to fluctuations?

partons or hadrons?

- ▶ species independence for ⟨p_t⟩
- need particle-identified fluctuations



Summary: Partial Thermalization

$\langle p_t \rangle$, p_t and charge fluctuations

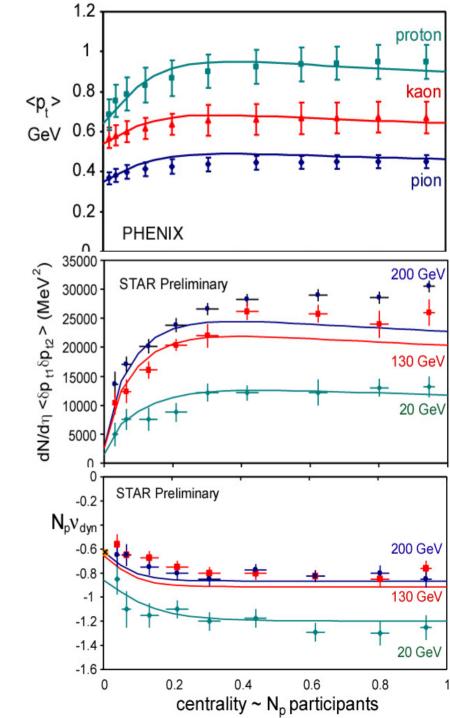
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Initial Fluctuations

AA collision:

M participant nucleons

 \rightarrow independent strings

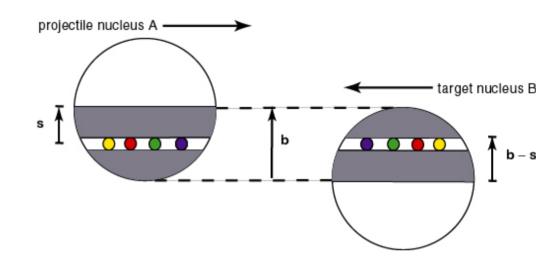
multiplicity $\propto M$

variance

$$R_{AA} \propto M^{-1}$$

dynamic p_t fluctuations

$$\left\langle \delta p_{t1} \delta p_{t2} \right\rangle \equiv \frac{1}{\left\langle N(N-1) \right\rangle} \left\langle \sum_{i \neq j} \delta p_{ti} \delta p_{tj} \right\rangle$$
$$\left\langle N(N-1) \right\rangle = \left\langle N \right\rangle^2 \left(1 + R_{AA}\right)$$



$$\left\langle \delta p_{t1} \delta p_{t2} \right\rangle \approx \frac{2 \left\langle \delta p_{t1} \delta p_{t2} \right\rangle_{pp}}{M} \left(\frac{1 + R_{pp}}{1 + R_{AA}} \right)$$

Fluctuations Near Equilibrium

correlations from non-uniformity – more likely to find particles near "hot spots" → spatial correlation function

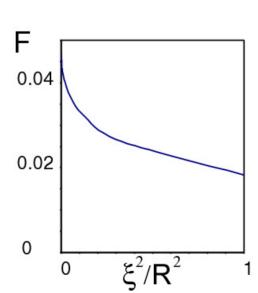
$$\langle \delta p_{t1} \delta p_{t2} \rangle \propto \iint \delta \overline{p}_{t1}(x_1) \delta \overline{p}_{t2}(x_2) r(x_1, x_2) dx_1 dx_2$$

assume:

- longitudinal Bjorken expansion
- uncorrelated longitudinal and transverse d.o.f.'s
- $\delta p_t \propto \delta T(x)$ independent of rapidity
- gaussian densities, correlation function

$$\langle \delta p_{t1} \delta p_{t2} \rangle_e \approx \frac{\langle p_t \rangle^2 R_{AA}}{1 + R_{AA}} F(\xi / R_t)$$

transverse size R_t , correlation length ξ



Approach to Equilibrium

Boltzmann Equation for phase space density f(x,p)

$$\frac{\partial f}{\partial t} + \vec{v}_{\vec{p}} \cdot \vec{\nabla} f = -\nu (f - f^e)$$

approximation: relaxation time v^{-1}

Bjorken scaling:

$$f = f_0(p_z', p_t, s)S(t, t_0) + \int_{t_0}^t f^e(p_z', p_t, s)S(t, s)vds$$

Baym; Gavin; B. Zhang & Gyulassy

where $p_z' = p_z(t/t_0)$

survival probability – chance of no scattering

$$S(t,t_0) = e^{-\int_{t_0}^t v \, ds}$$

$$\langle p_t \rangle = \langle p_t \rangle_0 S + \langle p_t \rangle_e (1 - S)$$

Nonequilibrium Fluctuations

Introduce fluctuations:

- Boltzmann Equation plus Langevin noise
- Langevin noise for fe

relaxation equation for correlation functions

$$P(x_1, p_1, x_2, p_2) = \langle f_1 f_2 \rangle - \langle f_1 \rangle \langle f_2 \rangle - \delta_{12} \langle f_1 \rangle$$

$$\mathbf{C}(x_1, p_1, x_2, p_2) = \left\langle f_1 f_2^e \right\rangle - \left\langle f_1 \right\rangle \left\langle f_2^e \right\rangle$$

$$\frac{dP_{12}}{dt} = -2\nu P_{12} + \nu(C_{12} + C_{21}); \qquad \frac{dC_{12}}{dt} = -\nu(C_{12} - P_{21}^{e})$$

$$\frac{d}{dt} = \frac{\partial}{\partial t} + \nu_{p_1} \cdot \nabla_1 + \nu_{p_2} \cdot \nabla_2$$

Bjorken flow:

$$P_{12} = P_{12}^{0} S^{2} + 2 C_{2}^{0} S (1 - S) + P_{12}^{e} (1 - S)^{2}$$

single particle:

$$f(p) = e^{-(E - \vec{p} \cdot \vec{v})/T}, \quad \vec{v} << 1$$

$$P^{e}(x_1, p_1, x_2, p_2) = f(p_1)f(p_2)A$$

$$A(x_1, p_1, x_2, p_2) = r(x_1, x_2) \leftarrow$$

$$+ s(x_1, x_2) (\vec{v}(x_1) \cdot \vec{p}_1) (\vec{v}(x_1) \cdot \vec{p}_1)$$

$$+ t(x_1, x_2) E_1 E_2$$
density correlations

single particle:

$$f(p) = e^{-(E - \vec{p} \cdot \vec{v})/T}, \quad \vec{v} << 1$$

$$P^{e}(x_1, p_1, x_2, p_2) = f(p_1)f(p_2)A$$

$$A(x_{1}, p_{1}, x_{2}, p_{2}) = r(x_{1}, x_{2})$$

$$+ s(x_{1}, x_{2})(\vec{v}(x_{1}) \cdot \vec{p}_{1})(\vec{v}(x_{1}) \cdot \vec{p}_{1})$$

$$+ t(x_{1}, x_{2})E_{1}E_{2}$$
momentum correlations

single particle:

$$f(p) = e^{-(E - \vec{p} \cdot \vec{v})/T}, \quad \vec{v} << 1$$

$$P^{e}(x_1, p_1, x_2, p_2) = f(p_1)f(p_2)A$$

$$\begin{split} A(x_1,p_1,x_2,p_2) &= r(x_1,x_2) \\ &+ s(x_1,x_2) \Big(\vec{v}(x_1) \cdot \vec{p}_1 \Big) \Big(\vec{v}(x_1) \cdot \vec{p}_1 \Big) \\ &+ t(x_1,x_2) E_1 E_2 & \qquad \qquad \text{temperature correlations} \end{split}$$

single particle:

$$f(p) = e^{-(E - \vec{p} \cdot \vec{v})/T}, \quad \vec{v} << 1$$

$$P^{e}(x_1, p_1, x_2, p_2) = f(p_1)f(p_2)A$$

$$A(x_1, p_1, x_2, p_2) = r(x_1, x_2)$$

$$+ s(x_1, x_2) (\vec{v}(x_1) \cdot \vec{p}_1) (\vec{v}(x_1) \cdot \vec{p}_1)$$

$$+ t(x_1, x_2) E_1 E_2$$