

FLOW AT RHIC :

WHAT SHOULD WE SEE AND HOW ?

J. Y. Ollitrault

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I FLOW AND THERMALIZATION

(is thermalization partial or complete?)

1. A dependence
2. Absolute magnitude of elliptic flow at RHIC
3. Flow of π , K , p : p_T dependence
4. Centrality dependence of elliptic flow

II METHODS FOR MEASURING THE FLOW

1. Existing methods and assumptions :
a brief historical survey
2. Minijets at RHIC
3. What should be done ?

VARIATION OF FLOW

WITH THE SIZE OF THE SYSTEM

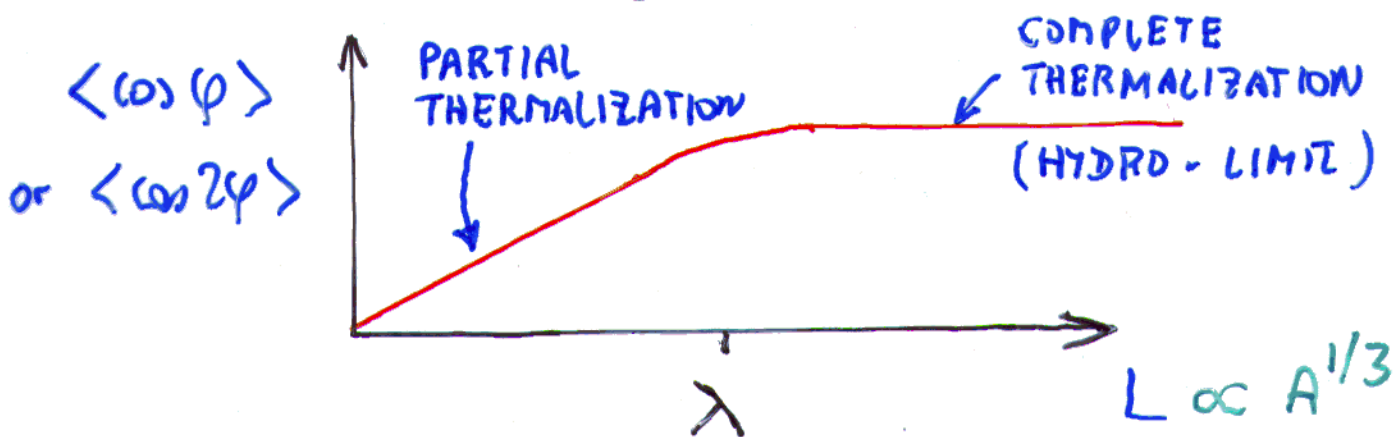
- FLOW AT HIGH ENERGIES (DIRECTED OR ELLIPTIC) RESULTS FROM COLLISIONS BETWEEN PARTICLES (NN, πN , $\pi\pi$...)
- LOCAL THERMAL EQUILIBRIUM (HYDRO-LIMIT) IS ACHIEVED IF

$$\lambda \ll L$$

MEAN FREE PATH
($= \frac{1}{\sigma n}$)

SIZE OF THE SYSTEM
($\propto A^{1/3}$)

- FLOW OBSERVABLES DEPEND ON L TYPICALLY LIKE (AT FIXED ENERGY)

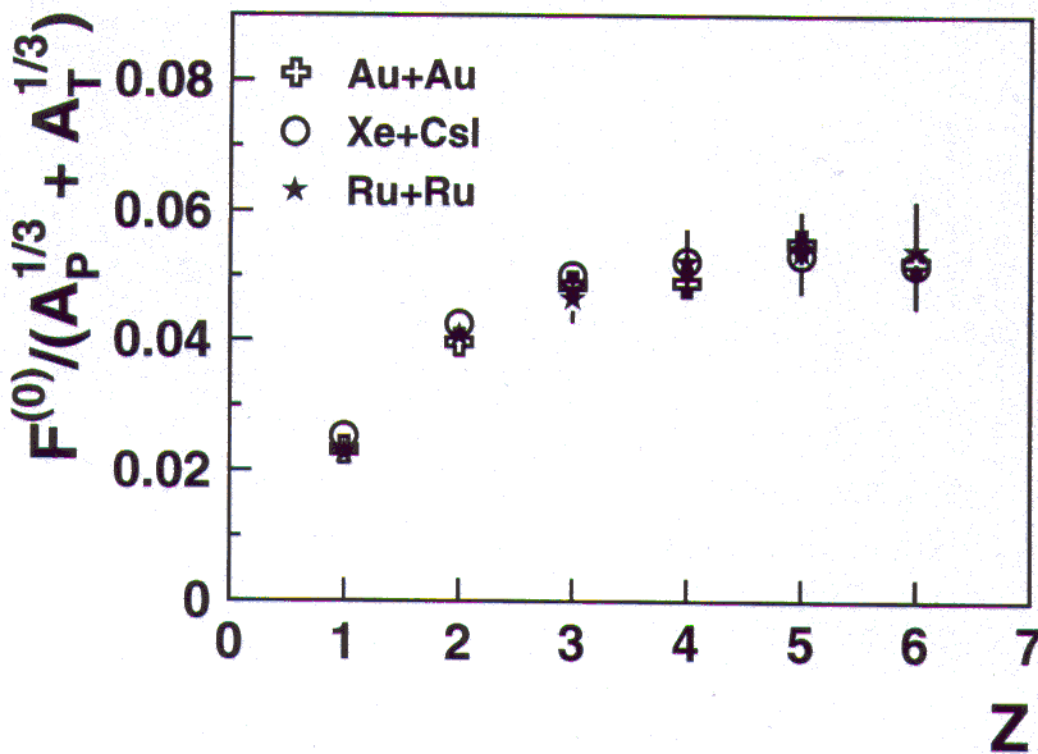


THE VARIATION OF FLOW WITH A GIVES A DIRECT INFORMATION ON THE DEGREE OF THERMALIZATION ACHIEVED IN THE SYSTEM

(SEE, HOWEVER, P. FILIP, Act. Phys. Slov 47 (1997) 53)
WHICH PREDICTS A DIFFERENT BEHAVIOUR

Figure 4

A - DEPENDENCE OF DIRECTED
FLOW AT $E/A \approx 400 \text{ MeV}$



FROM F. RAMI et al, FOPI Collaboration,
to be published - (Nov 1998)

DIRECTED FLOW SCALES LIKE $A^{1/3}$:
THERMALIZATION IS FAR FROM COMPLETE

(cf P. Danielewicz et al, PRC 38 (1988) 120)

ABSOLUTE MAGNITUDE OF IN-PLANE ELLIPTIC FLOW

THE A -DEPENDENCE OF DIRECTED AND ELLIPTIC FLOW HAS UNFORTUNATELY NOT BEEN STUDIED AT AGS AND SPS

HOWEVER, HYDRO MODELS FAIL TO PREDICT THE ABSOLUTE MAGNITUDE OF IN-PLANE ELLIPTIC FLOW OBSERVED AT SPS

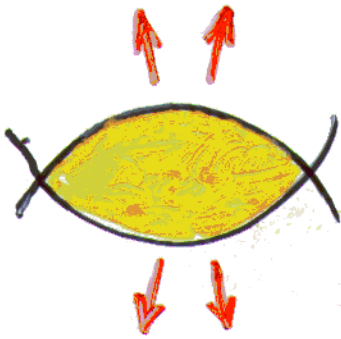
$$\begin{aligned} \text{SPS} &: v_2 \leq 4\% \text{ FOR } \pi, p \\ \text{HYDRO} &: v_2 \sim 20\% \end{aligned}$$

THIS SUGGESTS THAT ONLY PARTIAL THERMALIZATION IS ACHIEVED AT SPS

AT RHIC, THERMALIZATION SHOULD BE BETTER IF A LONG LIVED, STRONGLY INTERACTING PLASMA OF GLOUED QUARKS AND GLUONS IS FORMED.

TIME SCALE FOR

IN-PLANE ELLIPTIC FLOW



ELLIPTIC FLOW DEVELOPS
FOR $t \sim \frac{R}{c_s}$ (RADIUS)
VELOCITY OF SOUND

THROUGH IN-PLANE ELLIPTIC FLOW,
WE "SEE" THE SYSTEM AT $t \approx 10 \text{ fm}/c$,
I.E. AT A LATE STAGE ($\neq \psi, \chi, e^+e^- \dots$)

(A SHORT LIVED QUARK-GLUON PLASMA
WILL HAVE NO SIZEABLE EFFECT HERE)

RELEVANT DENSITY :

NOT THE "BJORKEN ESTIMATE" $\frac{1}{\pi R^2} \frac{1}{t_0} \frac{dN}{dy}$

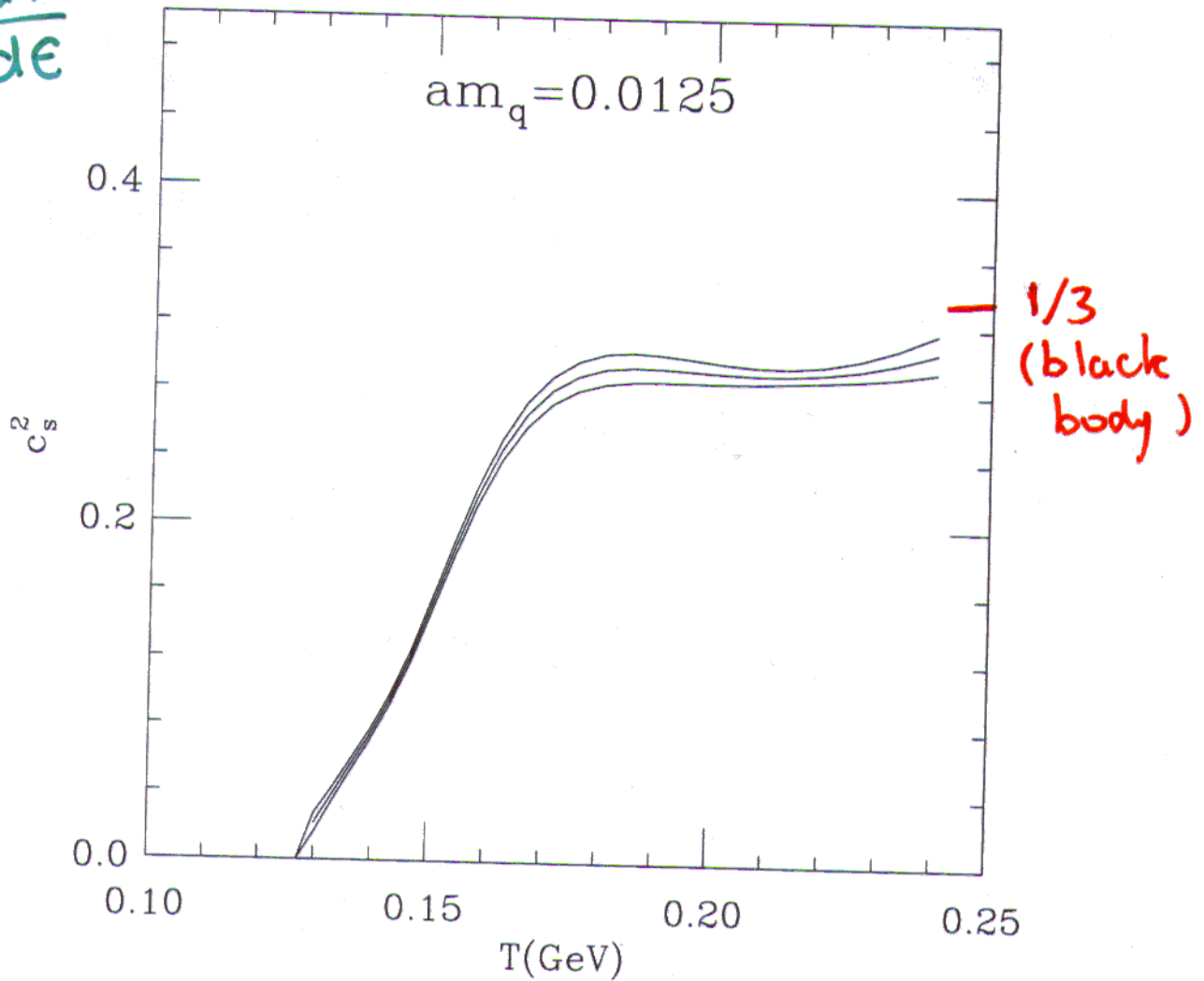
BUT RATHER ($t_0 \rightarrow R$)

$$\underline{n \sim \frac{1}{\pi R^3} \frac{dN}{dy}}$$

5-10 TIMES SMALLER
ROUGHLY INDEPENDENT OF A
(ONLY A SLIGHT INCREASE)

LATTICE CALCULATION OF THE SOUND VELOCITY

$$c_s^2 = \frac{dP}{dE}$$



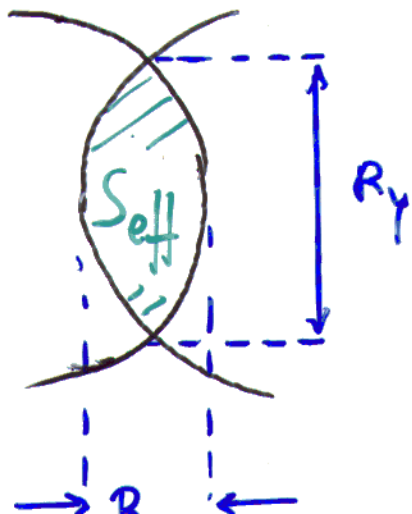
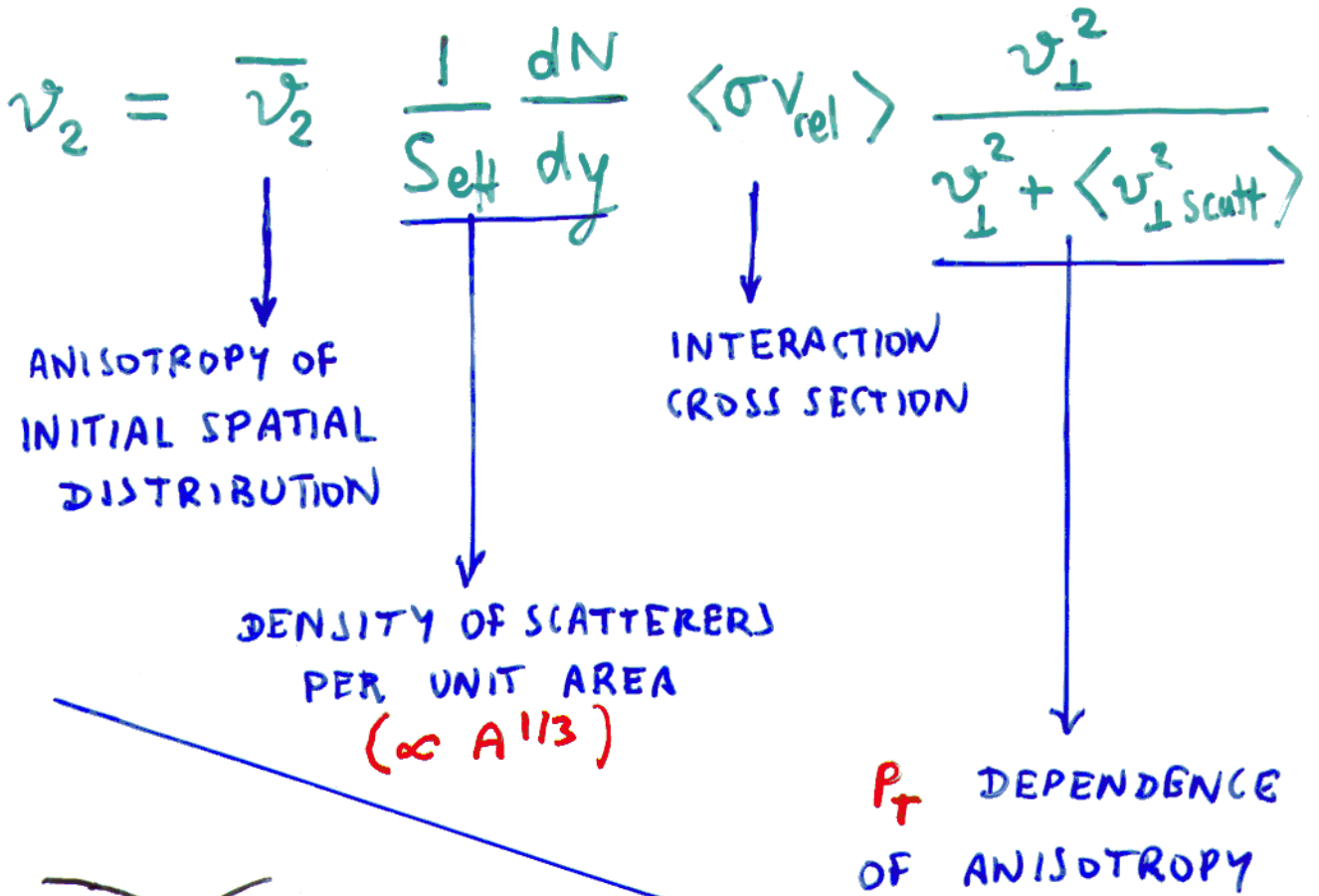
FROM MILC collaboration, hep-lat/9612025

→ ABOVE $T = 170$ MeV, THE VELOCITY OF SOUND IS CLOSE TO THE ULTRARELATIVISTIC LIMIT $c_s^2 = 1/3$

ELLIPTIC FLOW FROM PARTIAL THERMALIZATION

H. Heiselberg and A.M Levy, nucl-th/9812034

CALCULATE IN-PLANE ELLIPTIC FLOW IF
THERE IS **AT MOST ONE COLLISION PER
PARTICLE** :



$$\overline{v_2} = \frac{R_y^2 - R_x^2}{R_y^2 + R_x^2}$$

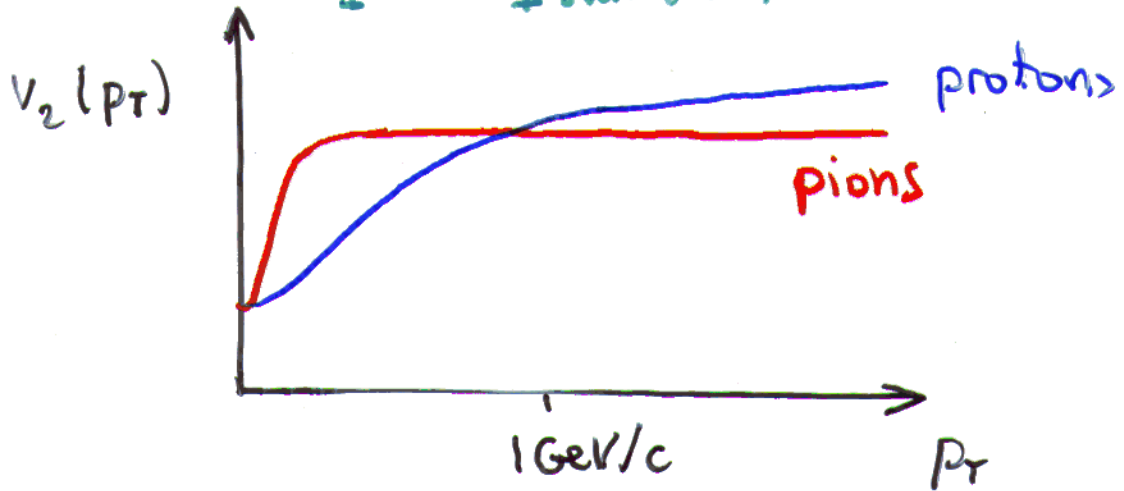
P_T DEPENDENCE OF ELLIPTIC FLOW :

PARTIAL VS COMPLETE THERMALIZATION

PARTIAL THERMALIZATION ($N_{\text{collisions/particle}} \ll 1$)

According to Heiselberg and Levy,

$$v_2(p_T) \propto \frac{v_{\perp}^2}{v_{\perp}^2 + \langle v_{\perp}^2 \text{ scatterer} \rangle}$$



→ p_T dependence depends on particle mass

→ $v_2(p_T)$ saturates at high p_T

COMPLETE THERMALIZATION ($N_{\text{collisions/particle}} \gg 1$)

of emitted particles $\propto e^{-p \cdot u / T}$

particle

4-momentum

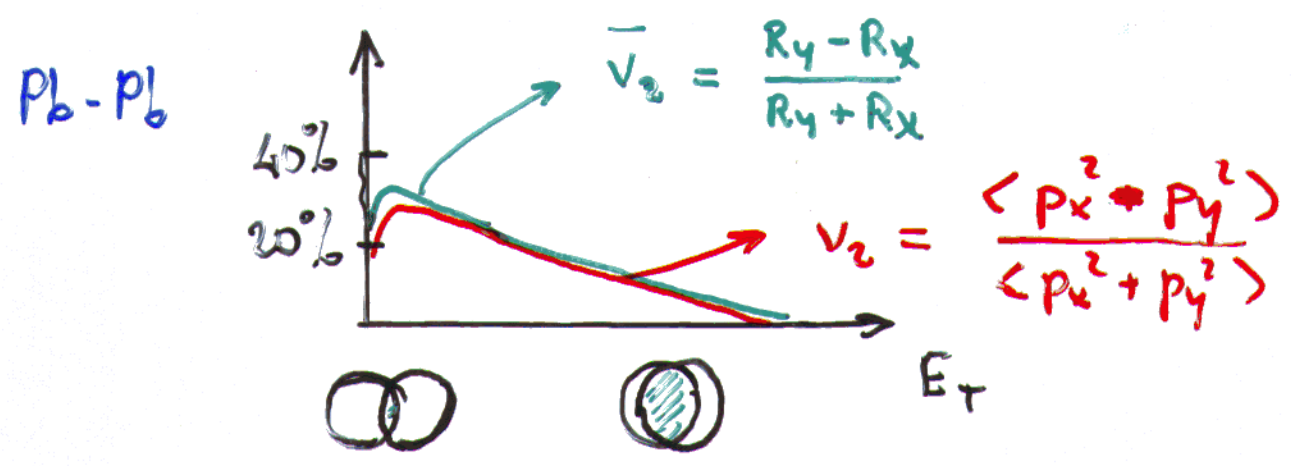
fluid

4-velocity

→ $v_2(p_T)$ same for \neq particles at given p_T
(S. Voloshin, PRC 55 (1997) 1630)

CENTRALITY DEPENDENCE OF ELLIPTIC FLOW

- THE ELLIPTIC FLOW CALCULATED IN HYDRO IS VERY CLOSE TO THE SPATIAL ANISOTROPY OF THE INITIAL DISTRIBUTION



THE MAXIMUM OCCURS FOR (IMPACT PARAMETER) LARGER THAN IN ACTUAL EXPERIMENT)

- IF ONLY PARTIAL THERMALIZATION OCCURS,

$$v_2 \propto \bar{v}_2 \underbrace{\frac{1}{\text{Self}} \frac{dN}{dy}}_{\text{density of scatterers / unit area}}$$

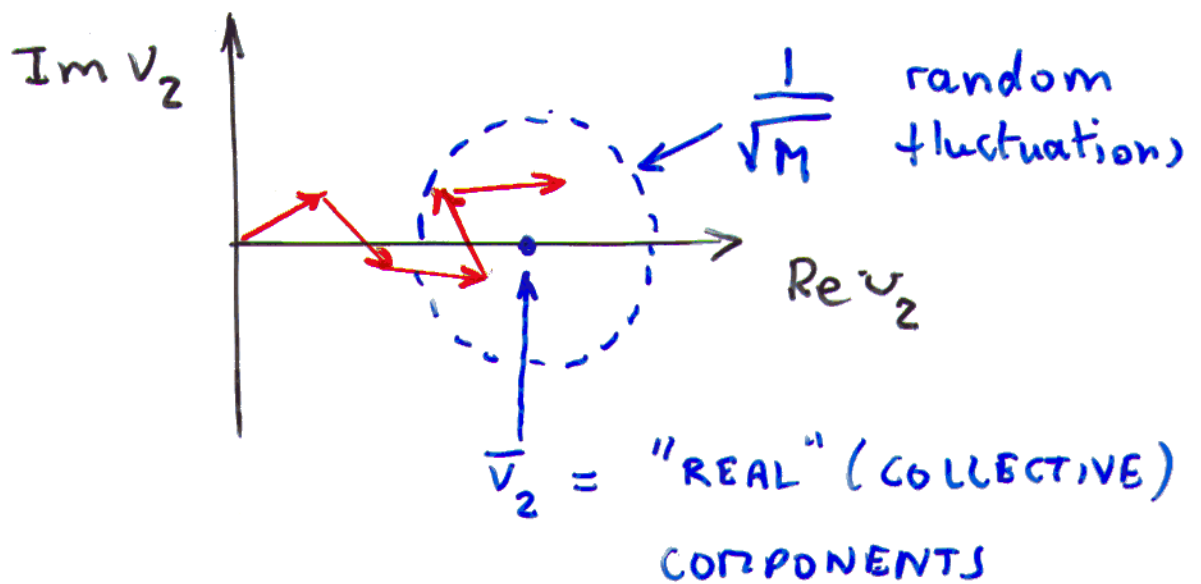
THE MAXIMUM IS SHIFTED TOWARDS LESS PERIPHERAL COLLISION)

(AND THE ABSOLUTE VALUE IS MUCH SMALLER)

METHODS FOR MEASURING THE FBW

THE GENERAL IDEA: USE A GLOBAL VARIABLE OBTAINED BY SUMMING OVER MANY PARTICLES SO THAT COLLECTIVE VELOCITIES SUM UP RANDOM (THERMAL) VELOCITIES CANCEL IN THE SUM

$$v_2 = \frac{1}{M} \sum_{j=1}^M e^{2i\varphi_j}$$



THE CRUCIAL QUANTITY IS THE RATIO OF COLLECTIVE AND STATISTICAL (RANDOM) COMPONENTS

$$\chi \approx \bar{v}_2 \sqrt{M}$$

FLOW IS OBSERVABLE UNAMBIGUOUSLY

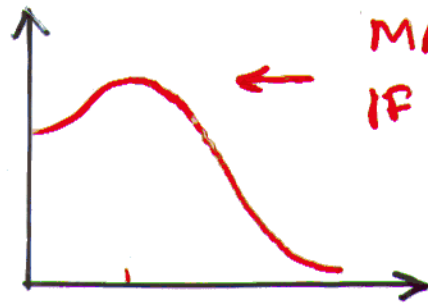
IF $\chi > 1$ (χ ranges from 0,3 to 2,4

TWO METHODS

Postkanzer, Voloshin, PRC58 (1998) 1671

1. PLOT THE DISTRIBUTION OF $|v_2|$

$$\frac{1}{|v_2|} \frac{dN}{d|v_2|}$$



MAXIMUM AT $|v_2| \neq 0$
IF $\chi > 1$

$|v_2|$

Danielewicz - Gyulassy PLB 129 (1987) 283

Ollitrault PRD46 (1992) 229

Barrette et al, E877, PRL 73 (1994) 2532

2. SUBEVENT METHOD :

Danielewicz, Odyniec, PLB 157 (1985) 146

MEASURE v_2 IN 2 RANDOMLY CHOSEN SUBEVENTS
AND STUDY THE ANGULAR CORRELATION BETWEEN
SUBEVENTS

EASIER (DEPENDS ONLY ON χ , DOES NOT REQUIRE
 $\chi > 1$)

BUT ASSUMES: ONLY CORRELATIONS ARE FROM FLOW:

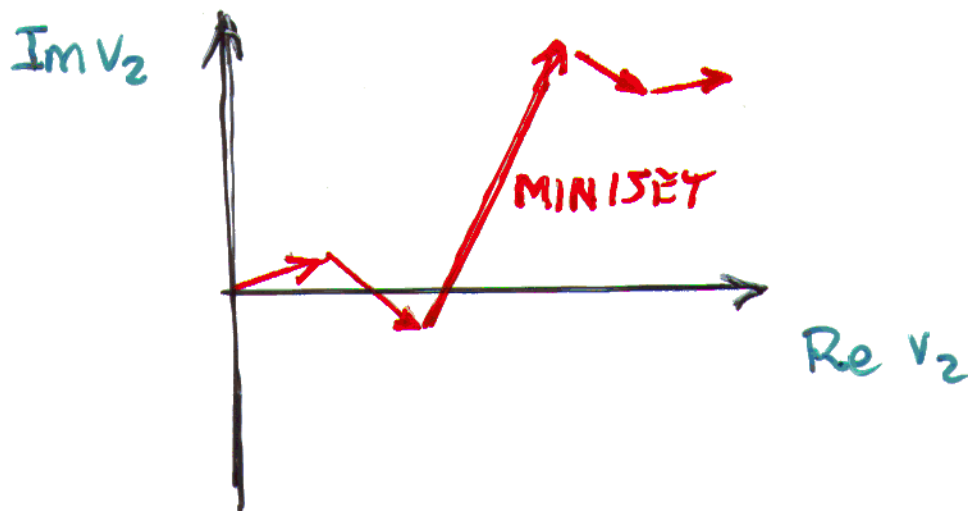
• UNCONTROLLABLE ASSUMPTION?

• SEEMS TO BE OK AT LOWER ENERGIES
(ONLY SIZEABLE CORRELATIONS ARE THOSE
FROM MOMENTUM CONSERVATION) BUT ...

MINIJETS, EVENT-BY-EVENT FLUCTUATIONS...

- MINIJETS, DCC's ... MAY PRODUCE STRONG AZIMUTHAL CORRELATIONS AMONG THE PRODUCED PARTICLES, WHICH HAVE A DIFFERENT ORIGIN \neq ELLIPTIC FLOW

- HOW DOES THIS AFFECT FLOW MEASUREMENTS?



→ INCREASE OF STATISTICAL, FINITE MULTIPLICITY FLUCTUATIONS WHICH MAY NOT SCALE LIKE $1/\sqrt{M}$

→ $\chi < v_2 \sqrt{M}$ MAY BE SMALLER THAN EXPECTED (ALTHOUGH BOTH v_2 and M ARE LARGER AT RHIC THAN AT SPS)

ARE FLUCTUATIONS GAUSSIAN?

- ASSUMPTION THAT FINITE MULTIPLICITY FLUCTUATIONS ARE GAUSSIAN ALLOWS TO PUSH METHOD 1 TOWARDS LOWER VALUES OF χ : STUDY DEVIATIONS FROM GAUSSIAN SHAPE.

- THE GAUSSIAN CHARACTER SHOULD BE TESTED FOR CENTRAL COLLISIONS (NO FLOW)

$$Q = \frac{1}{\sqrt{M}} \sum_{j=1}^M e^{2i\varphi_j}$$

↳ in order to scale by the fluctuations

CHECK THAT $\langle |Q|^4 \rangle = 2 \langle |Q|^2 \rangle^2$

IF FLOW IS PRESENT

$$\frac{\langle |Q|^4 \rangle}{\langle |Q|^2 \rangle^2} \approx 2 - \chi^4$$

FOR SMALL χ (DEVIATIONS FROM A GAUSSIAN ARE SMALL!)